

Superallowed beta decay

J.C. Hardy, I.S. Towner, V.E. Jacob, H.I. Park, N. Nica, M. Bencomo, T. Eronen,
V. Horvat, and L. Chen

Superallowed $0^+ \rightarrow 0^+$ beta decay between $T=1$ analogue states has been a subject of continuous and often intense study for five decades. The ft values of such transitions are nearly independent of nuclear-structure ambiguities and depend uniquely on the vector part of the weak interaction. Their measurement gives us access to clean tests of some of the fundamental precepts of weak-interaction theory, and, over the years, this strong motivation has led to very high precision being achieved both in the experiments and in the theory used to interpret them. We have a major program at the Cyclotron Institute to study superallowed beta decay.

To obtain the ft value for any transition, three quantities must be measured: the half-life $t_{1/2}$ of the parent, the Q_{EC} value for the transition of interest, and the branching ratio R for that transition. Our most recent complete survey of world data on these superallowed decays, published in 2015 [1], provides a critical evaluation of all the experimental data and final ft values obtained from the averaged results, to which radiative and isospin-symmetry-breaking corrections have been applied in order to derive a final set of “corrected ft values”, denoted $\mathcal{F}t$ for 14 transitions known to $\sim 0.1\%$ precision. Last year, we made an interim update of this survey [2], the results from which are shown in Fig. 1. Excellent consistency among the average $\mathcal{F}t$ values for all 14 transitions – an expected consequence of the conservation of vector current (CVC) – confirms the validity of the correction terms; and our recent measurement of ^{38}Ca decay [3,4], which closely compares a pair of mirror superallowed transitions with $A = 38$, further supports that validity.

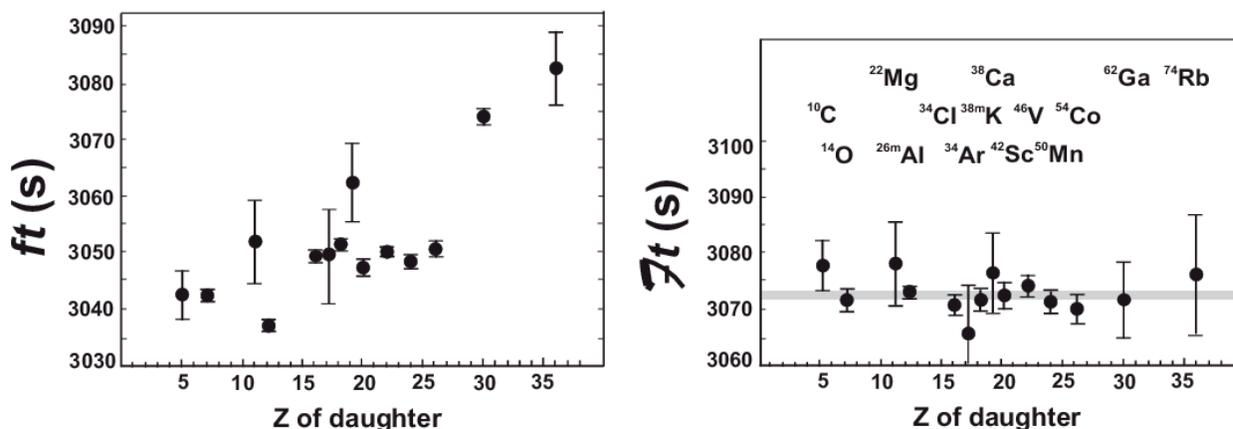


FIG. 1. Results from the 2015 survey [1] updated to 2016 [2]: The uncorrected ft values for the 14 best known superallowed decays appear on the left; the same results but incorporating the radiative and isospin-symmetry-breaking correction terms are on the right. The grey band in the right panel is the average $\mathcal{F}t$ value and its uncertainty.

The resultant average $\mathcal{F}t$ value, when combined with the muon lifetime, yields the up-down quark-mixing element of the Cabibbo-Kobayashi-Maskawa (CKM) matrix, $V_{ud} = 0.97420(21)$, a result

that is consistent with, but more precise than, values we have obtained in previous analyses of superallowed β decay. The unitarity test on the top row of the matrix becomes $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.99963$ (49) if the Particle Data Group recommended value for V_{us} is used. However, recent lattice QCD calculations have introduced some inconsistency into kaon-decay measurements of V_{us} and V_{us}/V_{ud} . In ref. [1], we have examined the impact of these new results on the unitarity test and conclude that there is no evidence of any statistically significant violation of unitarity. Finally, from the $\mathcal{F}t$ -value data we also set limits on the possible existence of scalar interactions.

This result is not only a significant verification of the standard model but the uncertainty quoted on the sum provides a tight limit on any possible new physics beyond the standard model, such as right-hand currents, extra Z bosons or supersymmetric models. In short, superallowed $0^+ \rightarrow 0^+$ beta decay provides a high-profile application of nuclear-physics measurements to the study of fundamental symmetries, a subject of vital interest to both nuclear and particle physicists. Although much has already been achieved in this field by nuclear physicists, improvements are still possible. Reducing the uncertainty on the unitarity sum – and, with it, the scope for new physics – remains the primary goal of our research program.

Our approach follows from the observation [1] that the second largest contributor to the uncertainty in V_{ud} is the theoretical uncertainty in the nuclear-structure-dependent corrections, δ_{NS} and δ_C , used in the derivation of the $\mathcal{F}t$ values. Though these corrections are only of order 1%, their effect is very significant: The bottom two panels of Fig. 1 show the result of applying the nuclear-structure-dependent corrections, δ_{NS} and δ_C (together with δ'_R , which is nearly independent of Z). Obviously they act very well to remove the considerable “scatter” in ft values apparent in the panel on the left, replacing it with the consistent set of corrected $\mathcal{F}t$ values appearing on the right. Since these corrections were determined [5] completely independently of the superallowed decay data, this consistency in $\mathcal{F}t$ values is already a powerful validation of the calculations, but obviously the remaining uncertainty still influences the final result for V_{ud} .

Even though the 2015 survey [1] included more than 222 individual measurements (and the 2016 update includes a dozen more) relating to 14 precisely known ft values, it is still possible for well selected experiments to make real improvements in the validation tests of the nuclear-structure-dependent correction terms. At TAMU we are currently focusing on adding to the ft -value list new superallowed transitions, selected from amongst those with *large* calculated corrections. If the ft values measured for cases with large calculated corrections also turn into corrected $\mathcal{F}t$ values that are consistent with the others, then this must verify the calculations' reliability for the existing cases, which have smaller corrections. We are studying decays from $T_z = -1$ parent nuclei, which consistently have higher predicted structure-dependent correction terms than the well-known $T_z = 0$ cases.

Of particular importance are the four $T_z = -1$ parent nuclei – ^{26}Si , ^{34}Ar , ^{38}Ca and ^{42}Ti – whose decays are mirrors to well-known superallowed decays from $T_z = 0$ parents. Specifically, the mirror-decay pairs are $^{26}\text{Si} \rightarrow ^{26m}\text{Al} \rightarrow ^{26}\text{Mg}$, $^{34}\text{Ar} \rightarrow ^{34}\text{Cl} \rightarrow ^{34}\text{S}$, $^{38}\text{Ca} \rightarrow ^{38m}\text{K} \rightarrow ^{38}\text{Ar}$ and $^{42}\text{Ti} \rightarrow ^{42}\text{Sc} \rightarrow ^{42}\text{Ca}$. Their importance stems from our observation that the ratio of mirror ft values for such cases is very sensitive to the model used to calculate the small isospin-symmetry-breaking corrections δ_{NS} and δ_C . The details have been described in our report on the first measurement of a mirror pair, with $A = 38$ [3]. Until very

recently, none of the $T_z = -1$ parent decays was known precisely enough to provide a statistically significant constraint on the correction terms via the ratio of mirror ft values, but we are now well on our way to rectifying this situation.

After a long period of incremental upgrades to our experimental techniques, we succeeded in pushing our precision in branching ratio measurements close to $\pm 0.1\%$, our ultimate goal. This is crucial for the characterization of $T_z = -1$ parent decays, which – unlike $T_z = 0$ decays – exhibit a number of strong Gamow-Teller branches that compete with the superallowed Fermi branch. A demonstration of our success in this endeavor is our measurement of the superallowed branching ratio for the decay of ^{38}Ca ($t_{1/2} = 444$ ms) to a precision of $\pm 0.2\%$, where that precision was actually limited by counting statistics, not systematics [3,4]. To our knowledge, this is the most precise direct branching-ratio measurement ever made for short-lived beta emitter. It also provides the first mirror pair of $0^+ \rightarrow 0^+$ superallowed emitters (^{38}Ca and $^{38}\text{K}^m$) that is precise enough to distinguish meaningfully between the Saxon-Woods-based radial-overlap correction, δ_{C2} , and the one based on Hartree-Fock radial wave functions. It favors the former over the latter, but we must await results from the other mirror pairs before we can be confident of the verdict.

We are now well embarked on the measurement of the remaining three accessible pairs. We have already completed measurements of the half-life and branching ratio for the superallowed decay of ^{34}Ar [6]. However, the branching-ratio result depends critically on the gamma-branching of the 666-keV level populated by beta decay in the daughter, ^{34}Cl : A possible weak branch from this level has 519 keV, which would be masked in our spectrum by the tail of the strong 511-keV annihilation peak. We are planning a (p, γ) measurement at Notre Dame to determine the relative intensity of this branch before we publish our results. In addition, we have made a successful measurement of the ^{26}Si beta-decay branching ratios, which is under analysis as the thesis project of M. Bencomo [7]. Finally we have made initial measurements of the half-life of ^{42}Ti , which have been analyzed [8]. A final measurement is scheduled for us to achieve the precision we seek.

We continue the endeavor to improve our data acquisition techniques for half-life measurements, making use of a TDC-based approach and a digital-pulse-analysis system, in addition to our longstanding analog technique. We are just concluding our analysis of a measurement of the half-life of another $T_z = -1$ superallowed emitter, ^{30}S , in which we used all three different methods for taking data [9]. In the same context, we continue to explore potential improvements in the statistical handling of our data [10].

With a somewhat different focus, in late 2015 we began a new measurement of the branching ratio for the superallowed decay of ^{10}C . Currently the uncertainty on the branching ratio dominates the uncertainty in the ^{10}C $\mathcal{F}t$ value. However, more interesting than just the precision of the $\mathcal{F}t$ value itself is its relationship to the world average of $\mathcal{F}t$ values for transitions in heavier nuclei, since the ^{10}C transition is the most sensitive to the possible presence of a scalar current. Currently the $\mathcal{F}t$ value for ^{10}C is slightly higher than the world average $\mathcal{F}t$ value, with an error bar that just about touches the world average value's error bar. If a more precise $\mathcal{F}t$ value of ^{10}C were found to deviate with greater statistical significance, it would be a signal for the existence of a scalar current. This work is still in progress [11].

Finally, each superallowed ft -value determination depends critically on the precision of the corresponding Q_{EC} -value measurement, which enters to the fifth power in the determination of f . In past

years we have made many such measurements with the JYFLTRAP Penning-trap mass spectrometer at the University of Jyväskylä cyclotron facility in Finland, where we collaborate with the team there. This facility is ideally suited to the measurement of Q_{EC} -values, which, with technical developments, have now reached sub-100-eV precision. So far we have measured the Q_{EC} -values for 10 superallowed emitters there: ^{10}C , $^{26}\text{Al}^m$, ^{34}Cl , ^{34}Ar , $^{38}\text{K}^m$, ^{38}Ca , ^{42}Sc , ^{46}V , ^{50}Mn , and ^{54}Co . It is now possible to improve some of the earliest measurements we made there. We have recently done so for the case of ^{42}Sc [12], reducing the uncertainty from ± 210 eV to ± 53 eV.

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